Irreversible Aging Dynamics and Generic Phase Behavior of Aqueous Suspensions of Laponite

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We study the aging behavior of aqueous suspension of Laponite having 2.8 wt % concentration using rheological tools. At various salt concentration all the samples demonstrate orientational order when observed using crossed polarizers. In rheological experiments we observe inherent irreversibility in the aging dynamics which forces the system not to rejuvenate to the same state in the shear melting experiment carried out at a later date since preparation. The extensive rheological experiments carried out as a function of time elapsed since preparation demonstrate the self-similar trend in the aging behavior irrespective of the concentration of salt. We observe that the exploration of the low-energy states as a function of aging time is only kinetically affected by the presence of salt. We estimate that the energy barrier to attain the low-energy states decreases linearly with increase in the concentration of salt. The observed superposition of all the elapsed time and the salt-concentration-dependent data suggests that the aging that occurs in low salt concentration systems over a very long period is qualitatively similar to the aging behavior observed in systems with high salt concentration over a shorter period.

I. Introduction

In soft glassy (nonergodic) materials, microstructure evolves to occupy low-energy sections of the energy landscape in its search for a possible thermodynamic equilibrium. The time required to attain a particular microstructural state, however, depends on the initial state and the nature of the path that connects these two states through an energy landscape. Greater degrees of freedom associated with the primary arrested entity (particle) of the glassy materials, such as nature of charges and charge distribution on the particle along with its shape and size distribution, may lead to a very intricate energy landscape. Consequently, the corresponding glassy material may show a very rich microstructural phase behavior. Aqueous suspensions of smectite clays are known to undergo ergodicity breaking and to form soft pasty materials that have significant applications in chemical, petroleum, cosmetic/pharmaceutical, food, etc., industries. Particularly, an aqueous suspension of Laponite is known to show a very rich phase behavior as a function of its concentration as well as the concentration of salt and is a subject of topical interest. However, the research work of more than a decade is divided over the possible microstates and their nomenclatures. Such contradiction is principally due to complexities involved in the charge distribution over the anisotropic particle of Laponite. In this article we have carried out extensive rheological study of aqueous suspension of Laponite having varying concentration of salt to investigate the effect of time elapsed since preparation on the phase behavior.

Laponite is a smectite hectorite clay and belongs to a structural family known as 2:1 phyllosilicates. The chemical formula of Laponite is Na$_x$[(Si$_8$Mg$_5$.5Li$_0$.3)O$_{20}$(OH)$_4$]$_{-0.7}$-$0.4$ in a unit cell of the Laponite crystal, six octahedral magnesium ions are sandwiched between two layers of four tetrahedral silicon atoms.

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These groups are balanced by 20 oxygen atoms and four hydroxyl groups. Isomorphous substitution of divalent magnesium by monovalent lithium induces deficiency of positive charge within the layer that gives a permanent negative charge to the surface of Laponite. The Laponite particle has a disklike shape with a diameter of 25−30 nm and thickness of 1 nm. In the powder form, Laponite particles are present in stacks with sodium ions in the interlayer gallery that balance the negative surface charge. In an aqueous medium, sodium ions dissociate rendering a net negative charge to the Laponite surface. The edge (rim) of the disk is composed of hydroxides such as Mg−OH and Si−OH. In the low-pH medium protonation of the edge is preferred rendering a −OH2⁺ (positive) charge to the edge. Intensity of the positive charge on the edge goes on decreasing with the increase in pH. In the limit of strong basic environment deprotonation dominates wherein H⁺ ions dissociate giving a −O⁻ (negative) charge to the edge. According to Tombácz and Szekeres, the point of zero charge for Si-OH group is associated with pH between 4.5 and 5.5, while according to Kosmulski, for MgOH the isoelectric point is expected to be above pH of 10. In a Laponite particle, the intermediate octahedral layer contains Mg (and Li in smaller amount), which forms complex with OH. In the outer tetrahedral layer Si has stronger association with oxygen on the surface of the Laponite. Therefore, the edge of the Laponite particle can be considered to have predominantly MgOH groups. As reported by Martin and co-workers, Laponite technical bulletin claims that the edge of Laponite particle contains predominantly MgOH and is positive below pH of 11. Tawari and co-workers estimated the edge charge using pH and conductivity measurements and found the same to be positive having magnitude 50e2 per Laponite particle at pH of 9.97. To avoid leaching of magnesium from the Laponite particle in an acidic environment, pH of the suspension in an aqueous media is generally kept around 10. In the absence of any externally added salt, there is a repulsion among the Laponite particles, which is believed to be causing ergodicity breaking at a very low concentration of salt. However, various experimental and simulation results suggest that the possibility of the attractive interactions between the edge and the surface cannot be ruled out. Addition of external salt such as NaCl increases the concentration of cations, which screen the negative charge on the surface and reduce the repulsion among the Laponite particles. Overall, the anisotropic disklike shape of the Laponite particle and the uneven charge distribution on the same along with effects arising from the concentration of Laponite and that of the externally added salt lead to a rich spectrum of microstructures, which evolve with time in a very complicated energy landscape.

In the literature, various groups have proposed different versions of the phase diagrams as a function of concentration of Laponite and that of the salt (NaCl). It is generally observed that soon after mixing around 1 vol % or more Laponite in water a soft solid state is formed that withstands its own weight. The microstructure of Laponite suspension is known to be very sensitive to the sample preparation protocol. At very low ionic concentration, ergodicity breaking is proposed to be caused by dominating repulsive interactions among the Laponite particles due to overlapping of double layers, and the corresponding state is identified as a repulsive glass. Tanaka et al. proposed that with increase in the concentration of salt, Coulomb repulsion among the particles leads to formation of an attractive glass. With further increase in the concentration of salt, van der Waals attraction prevails and the system enters a gel state. Some groups proposed that the ergodicity breaking in aqueous Laponite suspension is due to positive edge−negative surface attractive interactions which leads to the gel formation. Cordier and co-workers studied evolutions of elastic and viscous modulus for Laponite suspensions having varying concentration salt. They observed that the evolution of both the moduli shifts to lower ages for the systems having higher concentration of salt, suggesting that concentration of salt affects the rate of gelation. Ruzicka et al. proposed that at low concentration of salt and at low concentrations of Laponite the system forms an inhomogeneous structure which they represent as gel, while at low concentrations of salt and high concentrations of Laponite they observed homogeneous structure, which they represent as attractive glass. Li et al. suggested that the Laponite platelets interact via an attractive potential on short distances but repulsive on longer distances emphasizing the role of short-range attraction for the process of aggregation at higher volume fractions of Laponite. Recently, Jabbbari-Farouji and co-workers proposed that the repulsive glassy state obtained in the Laponite suspension having high Laponite concentration, and no added salt is similar to that observed in hard-sphere glasses. For the moderate concentrations of Laponite, they proposed that the gel state and the glass state are part of the same energy landscape, and the system can get arrested in the either state. Overall, there are multiple proposals, and over 15 years of research in this area has not lead to much consensus to unify various observations.

In this paper, we investigate evolution of elastic and viscous modulus of the aqueous Laponite suspensions having different concentrations of salt as a function of aging time at regular intervals for more than 18 days of idle time after the preparation of the sample. We observe that the suspensions undergo irreversible aging with subsequent evolution curves shifting to lower ages for the experiments carried out on a progressively later date since preparation. Experiments with various salt concentrations show self-similar behavior leading to salt concentration−idle time−aging time superposition.

II. Experimental Procedure

Laponite RD, procured from Southern Clay Products Inc., was dried in the oven at 120 °C for 4 h. The desired amount clay was mixed with deionized water, which has pH 10 and a predetermined amount of salt (NaCl), using an Ultra Turrex drive for a
that the third harmonic contributed negligibly to the response. 

always harmonic irrespective of the age. It was usually observed that the oscillatory shear stress having magnitude 10 Pa and frequency 0.1 Hz was applied in an oscillatory test with shear stress amplitude of 70 Pa and a frequency of 0.1 Hz. The time at which shear melting was stopped marked the beginning of the aging experiment ($t_i = 0$), as shown in Figure 1. In the subsequent aging experiments, oscillatory shear stress of the material.

In this work, the rheological experiments were carried out using a stress-controlled rheometer AR-1000 (Couette geometry with a bob diameter of 28 mm and a gap of 1 mm). The free surface of the suspension was covered with a thin layer of low viscosity silicon oil to avoid any contact of the same with air. We also measured the ionic conductivity of the suspensions using a CyberComm 6000 (Thermo Fisher Scientific Inc.) conductivity meter. All the experiments mentioned in this article were carried out at 10 °C.

III. Results

As shown in Figure 1, in each experiment, the shear melting step preceded the aging step. The inset in Figure 2 describes a shear melting behavior for a Laponite suspension having 3 mM NaCl concentration as a function of idle time ($t_i$). The experimental procedure is schematically described in Figure 1. After preparing the suspension, we carried out the rheological experiments at the desired idle times ranging from 3 days up to minimum 18 days. We also checked for the chemical stability of the suspension by carrying out the complexometric titrations using EDTA with Eriochrome black T indicator. We did not find any measurable amount of Mg$^{2+}$ ions in the suspension demonstrating the same to be chemically stable.

In the rheological study, in each experiment a fresh sample at certain idle time was loaded in a Couette cell using an injection syringe. The suspension partly melts in this step. After the thermal equilibrium was attained the sample was completely shear melted in an oscillatory test with shear stress amplitude of 70 Pa and a frequency of 0.1 Hz. The time at which shear melting was stopped marked the beginning of the aging experiment ($t_i = 0$), as shown in Figure 1. In the subsequent aging experiments, oscillatory shear stress having magnitude 10 Pa and frequency 0.1 Hz was applied to the shear melted suspensions to record the evolution of their viscoelastic properties. It should be noted that aqueous suspension of Laponite continuously undergoes structural evolution, increasing its viscosity and the elasticity causing broadening of the rheological linear response regime as a function of time. However, at very low age, since time required for the stress sweep experiment is usually more than the age of the sample, the suspension does not allow estimation of linear response regime. While the aging experiments were indeed in the linear response regime at higher ages, we confirmed that the oscillatory response of strain was always harmonic irrespective of the age. It was usually observed that the third harmonic contributed negligibly to the response.

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Figure 2. Plateau values of the complex viscosity in a shear melting step as a function of idle time ($t_i$) for various salt concentrations (filled symbols; squares: 0.1 mM; circles: 1 mM; up triangles: 3 mM; down triangles: 5 mM; stars: 7 mM). The inset shows day wise shear melting data for the 3 mM Laponite suspension. From bottom to top: $t_i = 3, 6, 9, 12, 15, 18,$ and 21 days. Rejuvenation was carried out at stress amplitude of 70 Pa.

Figure 3. Evolution of Storage ($G’$; filled symbols) and loss modulus ($G''$; open symbols) [stress amplitude 10 Pa, frequency 0.1 Hz] for a 5 mM suspension on $t_i = 3$ days for different rejuvenating stress amplitudes applied during the shear melting.

increases with increase in the idle time. This behavior demonstrates the partial irreversibility in the aging dynamics of aqueous suspension of Laponite, which causes complex viscosity in the shear melting step not to reach the same plateau value. It must be mentioned that, in the shear melting step we verified that the sample indeed underwent the complete rejuvenation. Such verification can be carried out by applying higher magnitude of shear stress amplitude to the sample in the shear melting step and studying the subsequent evolution of elastic and viscous modulus. Figure 3 shows evolution of elastic and viscous modulus after carrying out the shear melting at three different stresses for 5 mM Laponite suspension on $t_i = 3$ days. We observed that in the shear melting step the plateau value of complex viscosity was lower for greater stress. It can be seen that irrespective of the value of the rejuvenation stress in the mentioned range the subsequent evolution of both the moduli is not affected. This observation is in agreement with Joshi et al., who showed that evolution of the elastic modulus shows identical behavior irrespective of the amplitude of stress in the shear melting step when the shear melting is performed at the stress amplitudes greater than the yield stress of the material.

Subsequent to the shear melting experiments, the aging experiments were performed as shown in Figure 1. Parts a and b of Figure 4 show respective evolutions of $G’$ and $G''$ as a function of aging time ($t_a$) at various idle times ($t_i$) up to 52 days, for a system with 1 mM salt concentration. It can be seen that the evolutions of both $G’$ and $G''$ with respect to aging time gets systematically shifted to low aging times for the experiments carried out at a later date since preparation (higher idle time, $t_i$). Similar curvatures of

all the aging curves result in a natural superposition by horizontal shifting as described by the superposing gray points on $t_i = 18$ days data. We have also plotted the superposition of $G''$, which is shown in Figure 4b, in Figure 4a. It can be seen that $G''$ shows a two-step evolution with an initial increase followed by a decrease with the aging time. Correspondingly, $G'$ also shows a two-step evolution, wherein it undergoes a rapid rise in the first step, and after crossing over $G''$ around its maximum, it undergoes a slow power law increase in the second step. The crossover of $G'$ and $G''$ in an aging experiment is generally considered as an ergodicity breaking point, when the system shows a transition from a liquid-like to a solid (glass)-like behavior. Interestingly, for the experiment carried out for idle time $t_i = 52$ days, the system does not show the liquid regime ($G'' < G'$) and directly enters the glassy regime ($G' > G''$).

We carried out a similar superposition procedure for the evolution of $G'$ and $G''$ at various idle times for five salt concentrations in the range $0.1-7$ mM. Figure 5a,b shows respective superpositions of $G'$ and $G''$ obtained by shifting all the idle time-dependent data on to the data obtained on 18th day ($t_i = 18$ days). The same horizontal shift factors are required to obtain superposition of $G'$ as well as $G''$. However, in order to get superposition of $G''$, vertical shifting is needed in addition to the horizontal shifting. Interestingly, there is a remarkable similarity in Figures 4 and 5. The former figure represents the evolution of both the moduli at different idle times for 1 mM salt concentration, while Figure 5 represents the evolution of both the moduli for various salt concentrations on the idle time $t_i = 18$ days. Interestingly, similar to the behavior observed for the idle time dependence, the evolution of $G'$ and $G''$ also get shifted to the lower age with the increase in the concentration of salt. Additional similarities include two-stage increase in $G'$ with transition marked by maximum in $G''$, which also approximately corresponds to crossover of $G'$ and $G''$ (the aging time at which $G'$ shows maximum is $\sim 1.5$ times larger than the aging time at which $G''$ crosses over $G''$; however, these two points appear very close due to the logarithmic nature of the scale). For 7 mM salt concentration system, irrespective of the idle time, the value of $G'$ in the beginning of the experiment is always greater than $G''$ and shows only a second stage of the evolution. Likewise, the corresponding $G''$ also shows only a decrease with the aging time. Importantly, it should be noted that for $C_s = 1$ mM and on $t_i = 52$ days, the observed evolution of $G'$ and $G''$ shown in Figure 4 is qualitatively similar to the evolution data for 7 mM salt concentration shown in Figure 5.

We have plotted the horizontal and vertical shift factors shown in Figure 5a,b as a function of idle time ($t_i$) in Figure 6a,b, respectively. In general, the horizontal shift factors, which are common to both the moduli, increase with the idle time ($t_i$). This behavior suggests that the evolution of both the moduli at high idle times shifts to the low aging times irrespective of the concentration of salt. The vertical shift factor is needed only for the viscous modulus ($G''$) data to obtain the superposition. As shown in Figure 6b, the vertical shift factors remain almost close to unity for 0.1, 1, and 3 mM systems. However, for 5 and 7 mM salt concentrations, for which most of the evolution is in the nonergodic regime, the vertical shift factors decrease with the idle time.

We also performed the aging experiments for above-mentioned five salt concentrations at various frequencies in the range 0.1 to 10 Hz. It is observed that, for all the explored salt concentrations, evolution of $G'$ becomes independent of frequency after the crossover of $G'$ and $G''$. According to Fielding et al., this observation of weak dependence on frequency after the crossover of $G'$ and $G''$ again confirms the system to be in the nonergodic regime.


state. This weak frequency dependence implies that the shift factors shown in Figure 6a,b do not depend on frequency for the part of the superposition in the nonergodic regime.

The superpositions of the idle-time-dependent data for the five salt concentration systems, shown in Figure 5a,b, demonstrate that the evolutions of both the moduli for the systems having higher concentration of salt get shifted to a lower age preserving the overall curvature. In Figure 7 we have plotted the master superposition by horizontally and vertically shifting the concentration-dependent $G'$ and $G''$ superpositions plotted in Figure 5a,b. The respective superpositions at various salt concentrations are shifted to the superposition associated with 5 mM system, which results in “idle time—aging time—salt concentration superposition”. The corresponding horizontal shift factor, which is common for both the moduli, is plotted as a function of concentration ($C_s$) in Figure 8. It can be seen that the horizontal shift factor increases exponentially with increase in the concentration of salt ($C_s$). This master superposition and the associated trend of the horizontal shift factor implies that the aging that occurs in the

Figure 6. (a) Horizontal shift factor as a function of idle time ($t_i$) for storage and loss moduli. (b) Vertical shift factor as a function of idle time ($t_i$) for loss modulus (squares: 0.1 mM; circles: 1 mM; up triangles: 3 mM; down triangles 5 mM; stars: 7 mM).

Figure 7. A complete master curve combining both idle time ($t_i$) and salt concentration ($C_s$) dependence for representing evolution of (a) storage (filled symbols) and (b) loss modulus. (open symbols).

Figure 8. Horizontal and vertical shift factors used for obtaining the overall master curve for various salt concentration. Horizontal shift factor, $H$ (filled circles), which is common to both $G'$ and $G''$, shows exponential dependence on salt concentration. The vertical shift factor $V'$ is represented by stars while $V''$ is represented by open triangles.

In this section we will discuss the implications of the experimental results presented in the previous section with respect to the microstructure and the aging behavior of aqueous suspension of Laponite. As mentioned before, the microstructure of the aqueous suspension of Laponite is influenced by its anisotropic shape, uneven charge distribution on the particle, pH of the medium, and the concentration of dissociated Na$^+$ ions (either counterions or ions due to externally dissolved NaCl). These variables lead to the possibility of three types of interparticle interactions among the Laponite particles. The first interaction is the electrostatic repulsion between the negatively charged surfaces due to overlapping double layers. The second proposed interaction is the attraction between the negatively charged surface and the positively charged edge. The third proposed interaction is the van der Waals attraction between the Laponite discs at higher concentrations of salts when the negative charge on the Laponite surface is shielded by cations. Although the suspensions prepared using any protocol will involve mainly the above-mentioned interparticle interactions, one protocol may lead the suspension to a certain section of the energy landscape, which may not be accessible through another protocol.

It is usually believed in the literature that for the systems with low concentration of cations overall interactions among the Laponite particles are primarily repulsive. Increase in the concentration of cations (addition of salt) tends to shield the repulsive interactions, thereby reducing the energy


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required to reach the a certain level of reduced scattering intensity decreases exponentially with increase in the concentration of salt demonstrating an exponential dependence of the kinetics on the salt concentration \( [\tau \sim 10^{-0.42C}] \). Remarkably, their coefficient 0.42 matches closely with the coefficient 0.37 observed in present study (Figure 8).

Laponite, due to its oblate shape, is expected to show a nematic order in the aqueous suspensions at higher concentrations. Lemaire et al.\(^{60}\) and Gabriel et al.\(^4\) observed an orientational order in aqueous suspension of Laponite B above 2 wt % concentration. It should be noted that unlike Laponite RD, a synthetic layered silicate, used in this work, Laponite B is a synthetic layered fluorosilicate and has diameter 40 nm. Isotropic to biaxial and biaxial to nematic transition concentrations in oblate systems with hard surface interactions can be estimated by using Onsager’s approach.\(^{61}\) If \( a \) is the aspect ratio (\(L/D\)), the volume fractions at which discotic system undergoes isotropic to biaxial and biaxial to nematic transition are given by \( \phi_{n-b} = 0.33a \) and \( \phi_{b-n} = 0.45a \), respectively.\(^{61}\) For Laponite B, \( a = 1/40 \) leads to \( \phi_{n-b} = 0.00825 \), which gives weight fraction \( <0.02 \). Remarkably, this is precisely the same weight fraction for which both the above-mentioned papers observe order in their Laponite B suspension. Lemaire et al.\(^{60}\) observed the nematic order parameter of 0.55 ± 0.05 for 3 wt % Laponite B suspension. In principle, the structure factor should be above 0.8 for such transition in hard disks.\(^{61}\) They attribute this low order parameter to the topological defects that remain in their oriented samples.\(^{60}\) For Laponite RD, various experimental studies have proposed 3−3.5 wt % to be the concentration range above which order exists in the Laponite particles.\(^{20,27}\) In order to investigate state of Laponite suspension studied in the present work, we studied the birefringence behavior of the aged samples by placing the same between two crossed polarizers. Interestingly, we observed birefringence in all the explored salt concentration systems, suggesting our aged 2.8 wt % system to have an orientational order. If we consider Laponite RD to have diameter of 30 nm (usually between 25 and 30 nm), Onsager’s approach gives weight fraction above which order should be observed to be 0.0275 (2.75 wt %). However, in the case of Laponite, an edge of the particle is positively charged while the surface has a negative charge. In addition, the system is out of equilibrium and physically jammed. Under such situations the Onsager’s approach, which is strictly based on hard surface interactions and proposes liquid-to-liquid crystalline nematic first-order thermodynamic transition, may not be applicable. Nonetheless, the close match of the concentrations at which the order is observed in Laponite B as well as Laponite RD with that of the scaling arguments of Onsager is striking.

In the literature various possible structural scenarios for the aqueous suspension of Laponite have been proposed as a function of concentration of Laponite and concentration of NaCl. Since the edge of a Laponite particle is positively charged while the face has a permanent negative charge, one of the proposals suggests that aging in this system involves the Laponite particles reorienting themselves slowly to form the edge-to-surface interactions as a function of time.\(^{18,36,44}\) It is believed that the positive edge−negative surface interaction between the Laponite particles is significantly strong.\(^{18}\) Therefore, there is a possibility that while rejuvenating the overall suspension, the high stress cannot break all the edge-to-surface interactions causing partial irreversibility.


\(^{(59)}\) The horizontal shift factor is expected to depend on the barrier height (\(U_{b} \)) as \(H \sim -\exp{−U_{b}/k_{B}T} \). The present relation (ln(\(H(C_{b}) \sim C_{b} \)) therefore suggests that the barrier height decreases linearly with \(C_{b} \).


With increase in the idle time, the number density of such interactions increases, enhancing the number of edge–surface interactions that cannot be broken by the high stress, causing a steady enhancement in the plateau viscosity of the suspension in the shear melting step. In the subsequent aging experiments, since the aging process begins at a more matured (low energy) state, for the experiments carried out at a greater idle time, the overall evolution gets shifted to a lower age. As mentioned before, increase in the concentration of NaCl reduces the repulsion among the Laponite particles. Therefore, according to edge-to-surface aggregation school (house of cards structure), with the decrease in the extent of repulsion among Laponite particles, the rate of formation of edge-to-surface interactions is also expected to increase. This should lead to a faster buildup of the structure for the experiments carried out at a higher concentration of salt, making the evolution of both the moduli shift to a lower age. At very high concentration of salt, the formation of the edge-to-surface interactions is so fast that system directly enters the nonergodic state (\(G' > G''\)) after the shear melting stops. Therefore, even though the time taken by the low salt concentration suspension to reach a certain modulus level is significantly large compared to the same for high salt concentration system, the overall evolution maintains the same curvature on a logarithmic time axis leading to the superposition shown in Figure 7, thereby explaining the observed rheological behavior.

While the house of cards proposal (edge-to-surface interactions) explains the observed irreversibility in the aging behavior in the rheological experiments carried out in this work, the observed anisotropy in the Laponite suspension poses a strong challenge to this proposal. In principle, the collapsed house of cards structure (where the contact between two Laponite disks is bent) may explain the observed order and lowering of the nematic order parameter. However, in order to form such structure with bent contacts that fills the whole space, enough number density of Laponite particles needed in the system. If we consider \(n\) to be number of Laponite particles per unit volume, the volume fraction (\(\phi_v\)) of Laponite in water, in the limit of low volume fraction of Laponite, is given by \(\phi_v = \frac{\pi R^2 h}{4} n\), where \(R\) is radius of Laponite disk (= 13–15 nm) and \(h\) is thickness of Laponite particle (≈1 nm). Furthermore, the volume fraction of the spheres encompassing the \(n\) particles is given by \(4\pi R^3 / 3 = 4\phi_v / (Rh)\). For a system having concentration of 2.8 wt% the volume fraction of Laponite in water is around 0.011, which leads to the volume fraction of the spheres encompassing the particles at this concentration to 0.22. This volume fraction is too low to form space filling structure with bent interparticle contacts necessary to demonstrate an order. Moreover, the average interparticle distance for disk like particles is given by \((\pi R^2 h / \phi_v)^{1/3}\). This leads to an average interparticle distance of around 40 nm for a system having concentration of 2.8 wt%, which is greater than the diameter of the particle. Therefore, this analysis suggests that collapsed house of cards structure with bent angle of contact may not form a space filling arrested system at the concentration of 2.8 wt%.

To summarize discussion in this section, it is apparent that the length scale associated with repulsive interactions among the Laponite particles due to electrostatic screening is much smaller than the diameter of the particle for all the studied concentrations of salt. In addition, the edge has positive charge. The possible house of cards structure where there is a bond between positive edge and negative surface explains irreversible aging and the salt concentration dependence observed in the rheological experiments but does not explain the origin of anisotropy in the structure. Therefore, the possible microstructure of the system should be such that the particles remain in a long-range orientational order with interparticle interactions, the source of which is not yet clear to us, causing nonergodicity. The corresponding exploration of an energy landscape with respect to aging time by the system then occurs in such unique way that does not allow reversibility upon application of deformation field and demonstrates the kinetic dependence on the salt concentration. We feel that in order to get more insights into the present system the rheological and birefringence experiments performed in this work are needed to be carried out over a broader spectrum of concentrations of Laponite. In addition, there is clearly a need for more work on this topic particularly on realistic simulations to propose such microstructures as a function of system variables that will unify various observations of scattering experiments and rheological experiments mentioned in the literature.

**V. Conclusions**

In this work we study the aging behavior of aqueous suspension of Laponite having 2.8 wt % concentration as a function of the idle time (time elapsed since preparation of the suspension) and the concentration of NaCl. Interestingly, aged samples of Laponite suspension at all the studied salt concentrations show birefringence when observed using crossed polarizers, suggesting anisotropic nature of the microstructure of the same. In rheological experiments, we observe that the plateau value of the complex viscosity obtained in the shear melting experiments carried out at greater idle times increases for all the explored salt concentrations, thereby demonstrating partial irreversibility in aging. In addition, the subsequent evolution of the elastic and the viscous modulus, which represents aging, shifts to lower aging times for the experiments carried out at greater idle times. Interestingly, all the idle-time-dependent evolution data show the superposition suggesting time-separated self-similar aging behavior. In the salt-concentration-dependent experiments, evolution of elastic and viscous moduli shifts to lower aging times for higher concentration of salt, thereby showing kinetic dependence of rate of aging on the salt concentration. Remarkably, horizontal shifting of the salt concentration- and idle-time-dependent data shows the “idle time–aging time–salt concentration superposition”, suggesting self-similarity of the aging dynamics in aqueous suspension of Laponite. We estimate that the energy barrier to explore the low-energy states during aging process decreases linearly with increase in the concentration of salt. Remarkably, the idle-time- and the salt-concentration-dependent evolution of both the moduli superpose to show a generic trend. These results imply that the aging that occurs over a very long period in low salt concentration systems is qualitatively similar to that occurs in high salt concentration systems over a short period in the aqueous suspension of Laponite.

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